

The Chemical Evolution of Mg Isotopes vs. the Time Variation of the Fine Structure Constant

T. Ashenfelter and Grant J. Mathews
*Department of Physics, Center for Astrophysics,
University of Notre Dame, Notre Dame, IN, 46556*

Keith A. Olive
*William I. Fine Theoretical Physics Institute,
University of Minnesota, Minneapolis, MN 55455, USA*

The many-multiplet method applied to high redshift quasar absorption spectra has indicated a possible time variation of the fine structure constant. Alternatively, a constant value of α is consistent with the observational analysis if a non-solar isotopic ratio of $^{24,25,26}\text{Mg}$ occurs at high redshift. In particular, a higher abundance of the heavier isotopes $^{25,26}\text{Mg}$ are required to explain the observed multiplet splitting. We show that the synthesis of $^{25,26}\text{Mg}$ at the base of the convective envelope in low-metallicity asymptotic giant branch stars, combined with a simple model of galactic chemical evolution, can produce the required isotopic ratios and is supported by recent observations of high abundances of the neutron-rich Mg isotopes in metal-poor globular-cluster stars. We conclude that the present data based on high redshift quasar absorption spectra may be providing interesting information on the nucleosynthetic history of such systems, rather than a time variation of fundamental constants.

Over the last several years, there has been considerable excitement over the prospect that a time variation in the fine structure constant may have been observed in quasar absorption systems [1] - [4]. While many observations have led to interesting limits on the temporal variation of α (see [5] for a recent review), only the many-multiplet method [6] has led to a positive result, namely that $\frac{\delta\alpha}{\alpha} = (0.54 \pm 0.12) \times 10^{-5}$ over a redshift range of $0.5 < z < 3$, where $\delta\alpha$ is defined as the present value minus the past one.

The sources of systematic errors in this method have been well documented [3, 4] (see also [7]). Here, we would like to focus on one of these sources of systematic error for which there is recent evidence of a new interpretation, namely the isotopic abundances of Mg assumed in the analysis. The analyses in [1] - [4], have assumed terrestrial ratios for the three Mg isotopes. They have also shown that had they neglected the presence of the neutron rich Mg isotopes, the case for a varying α would only be strengthened. They further argued, based upon the galactic chemical evolution studies available at that time, that the ratio of $^{25,26}\text{Mg}/^{24}\text{Mg}$ is expected to decrease at low metallicity making their result a robust and conservative one.

In this paper, we will show that it is in fact quite plausible that the $^{25,26}\text{Mg}/^{24}\text{Mg}$ ratio was sufficiently *higher* at low metallicity to account for the apparent variation in α . As such, we would argue that while the many-multiplet method of analysis does not unambiguously indicate a time variation in the fine structure constant, it can lead to important new insights with regard to the nucleosynthetic history of quasar absorption systems.

Regarding the observations of Mg isotopes, Gay and Lambert [8] determined the Mg isotopic ratios in 20 stars in the metallicity range $-1.8 < [\text{Fe}/\text{H}] < 0.0$ with the aim of testing theoretical predictions [9]. (The notation $[\text{A}/\text{B}]$

refers to the the log of the abundance ratio A/B *relative* to the solar ratio.) Their results confirmed that the $^{25,26}\text{Mg}$ abundances relative to ^{24}Mg appear to decrease at low metallicity for normal stars. Although many stars were found to have abundance ratios somewhat higher than predicted, even the ‘peculiar’ stars which show enrichments in $^{25,26}\text{Mg}$ do not have abundance ratios substantially above solar.

Recently, however, a new study of Mg isotopic abundance in stars in the globular cluster NGC 6752 has been performed [10]. This study looked at 20 bright red giants which are all at a relatively low metallicity adopted at $[\text{Fe}/\text{H}] = -1.62$. These observations show a considerable spread in the Mg isotopic ratios which range from $^{24}\text{Mg}:^{25}\text{Mg}:^{26}\text{Mg} = 84:8:8$ (slightly poor in the heavies) to $53:9:39$ (greatly enriched in ^{26}Mg). The terrestrial value is $^{24}\text{Mg}:^{25}\text{Mg}:^{26}\text{Mg} = 79:10:11$ [11]. Of the 20 stars observed, 15 of them show ^{24}Mg fractions of 78% or less (that is below solar), and 7 of them show fractions of 70% or less with 4 of them in the range 53 - 67 %. This latter range is low enough to have a substantial effect on a determination of α in quasar absorption systems if the same ratios were to be found there. A previous study [12] also found unusually high abundances of the heavy Mg isotopes in M13 globular-cluster giants. Ratios of $^{24}\text{Mg}:^{25,26}\text{Mg}$ were found as low as 50:50 and even 44:56. Similar results were very recently found in [13]. According to [4], raising the heavy isotope concentration to $^{24}\text{Mg}:^{25,26}\text{Mg} = 63:37$ would sufficiently shift the multiplet wavelengths to eliminate the need for a varying fine structure constant. While dispersion in the data could be a symptom of systematic errors often occurring when data from several samples are combined, real dispersion is a signal that the observed abundances have been affected by local events.

Available calculations of Type II supernova yields in-

put into chemical evolution models [9], as well as observations of the Mg isotopic abundances in relatively low metallicity stars, support the idea that the heavy Mg isotopes were rarer in the past. Nevertheless, this conclusion is very sensitive to the star formation history in object under consideration. Intermediate mass stars in their giant phase, are expected to be efficient producers of $^{25,26}\text{Mg}$ [14, 15, 16] and the recent data from a low-metallicity globular cluster [10, 13] indicates substantial amounts of variation in the $^{25,26}\text{Mg}/^{24}\text{Mg}$ ratios as well as a number stars highly polluted in the neutron rich Mg isotopes. These observations show isotopic ratios considerably higher than the ratios predicted in zero metallicity supernovae and they conclude that asymptotic giant branch (AGB) stars may be responsible for this contamination.

Mg is produced in both Type I and Type II supernovae. In Type II supernovae, it is produced [17] in the carbon and neon burning shells with an abundance somewhat less than 10% of the oxygen abundance produced in massive stars. However, not much $^{25,26}\text{Mg}$ is produced in conventional stellar evolution at low metallicity. This is because the isotopes $^{25,26}\text{Mg}$ are produced primarily in the outer carbon layer by the reactions $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ and $^{25}\text{Mg}(n, \gamma)^{26}\text{Mg}$. According to the models of Woosley and Weaver [17], solar metallicity models produce final Mg isotope ratios reasonably close to solar values, while more massive stars tend to be slightly enhanced in the heavy isotopes, (e.g., the 25 M_{\odot} model gives a ratio of 65:15:20). Furthermore, the abundance of $^{25,26}\text{Mg}$ scales linearly with metallicity in the carbon shell. Hence, it would naively be expected that the ejecta from the first generation of supernovae would show a severe paucity of $^{25,26}\text{Mg}$.

Much of the solar abundance of Mg is produced in Type Ia supernovae but with Mg to Fe ratios below solar. For example, the models of Thielemann, Nomoto and Yokoi [18], give $[\text{Mg}/\text{Fe}] \simeq -1.2$. Due to the absence of free neutrons, essentially no $^{25,26}\text{Mg}$ are produced in Type Ia supernovae.

The results of chemical evolution models tracing the Mg isotopes were presented in [9]. These models clearly show the effect of the low yields of $^{25,26}\text{Mg}$ at low metallicity and predict $^{25,26}\text{Mg}/^{24}\text{Mg} < -0.8$ at $[\text{Fe}/\text{H}] < -1$. The result of [9] is essentially reproduced in the dashed curve shown in Figure 1 and discussed below.

Based on the conventional theory of Mg production (e.g. [17], models of galactic chemical evolution [9] and previously available data on the isotopic abundance of Mg [8], the adoption of solar isotopic Mg ratios by [1] - [4] in the many-multiplet analysis would appear to be a safe and conservative one. These models indicate that the $^{24}\text{Mg}:^{25,26}\text{Mg}$ ratio was higher than 79:21 in the past and it was shown [4] that a higher ratio only strengthens the case for a varying α . These models, however, do not include contributions from intermediate mass stars as we now describe. Recall that increasing the abundances of the heavier Mg isotopes would yield a larger value for

α , and a ratio of $^{24}\text{Mg}:^{25,26}\text{Mg} = 63:37$ is sufficient to obviate the need for a varying fine structure constant.

Recently, it has been appreciated [14, 15, 16] that intermediate mass stars of low metallicity can also be efficient producers of the heavy Mg isotopes during the thermal-pulsing AGB phase. Heavy magnesium isotopes (and to some extent silicon isotopes as well) are synthesized via two mechanisms both of which are particularly robust in 2.5-6 M_{\odot} stars with low metallicity. Such low-metallicity objects, indeed are precisely the kinds of objects which ought to produce the abundances observed in QSO narrow-line absorption systems at high redshift.

One process is that of hot-bottom burning. This process is also believed to be a copious source of lithium in low metallicity stars (cf. [19]) During the AGB phase, stars develop an extended outer convective envelope. Material in this convective envelope is mixed downward to regions of high temperature at the base. Of particular interest for this paper is that the base of the envelope is more compact and of higher temperature in low metallicity stars than in stars of solar composition. This can be traced to the decreased opacity of these objects. Furthermore, these stars would also have a shorter lifetime because they are hotter. Low to intermediate mass stars would contribute to the enrichment of the interstellar medium considerably sooner than their higher metallicity counterparts.

Because these stars become sufficiently hot ($T \geq 7 \times 10^7$ K), proton capture processes in the Mg-Al cycle become effective. Proton capture on ^{24}Mg then leads to the production of ^{25}Mg (from the decay of ^{25}Al) and to ^{26}Al (which decays to ^{26}Mg).

A second contributing process occurs deeper in the star during thermal pulses of the helium-burning shell. The helium shell experiences periodic thermonuclear run-aways when the ignition of the triple-alpha reaction occurs under electron-degenerate conditions. Due to electron degeneracy, the star is unable to expand and cool. Hence, the temperature rapidly rises until the onset of convection to transport the energy away. During these thermal pulses, ^{22}Ne is produced by α captures on ^{14}N which itself is left over from the CNO cycle. Heavy magnesium isotopes are then produced via the $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$ and $^{22}\text{Ne}(\alpha, \gamma)^{26}\text{Mg}$ reactions. It was argued recently [15], that in intermediate mass stars which experience a 3rd dredge-up, significantly greater amounts of $^{25,26}\text{Mg}$ are produced. A key point is that even though seed material is less plentiful in low metallicity stars, the reactions are very temperature sensitive. Hence, the increased temperature in the interior of low-metallicity stars more than compensates for the depleted seed material, leading to significant production of the heavy Mg isotopes. It has even been argued that these processes may also be net destroyers of ^{24}Mg [14, 16] due to the extreme temperatures attained.

To illustrate the effects of producing enhanced abundances of $^{25,26}\text{Mg}$ in intermediate mass stars, we show the results of a simple model of galactic chemical evo-

lution which traces the Mg isotopic abundances with and without the AGB source. When combined with the recent data showing such enhancements, our results suggest a plausible alternative for the interpretation of the quasar absorption system data based on the many-multiplet method.

For our purposes a simple recalculation of the results of Timmes et al. [9] with and without the contribution from intermediate-mass AGB stars is sufficient. This allows us to make a direct comparison with the conclusions of the previous authors.

The galactic chemical evolution model of Timmes et al. [9] is based upon exponential infall and a Schmidt star formation rate. We utilize a slightly modified model with updated yields [20] which nevertheless reproduces the results of [9] in the appropriate limit. Hence, we write the evolution of the surface density σ_i of an isotope i as,

$$\begin{aligned} \frac{d\sigma_i}{dt} = & \int_{0.8}^{40} B(t - \tau(m)) \Psi(m) X_i^S(t - \tau(m)) dm \\ & + \int_{2.5}^{9.0} B(t - \tau(m)) \Psi(m) X_i^{AGB}(t - \tau(m)) dm \\ & + m_{CO} X_i^{Ia} R_{Ia} - B(t) \frac{\sigma_i}{\sigma_{gas}} + \dot{\sigma}_{i,gas}, \end{aligned} \quad (1)$$

where $B(t)$ is the stellar birthrate at t , $\Psi(m)$ is the initial mass function (IMF), X_i^S is the mass fraction of isotope i ejected from single star evolution and Type II supernovae, and $\tau(m)$ is the lifetime of a star of mass m . The 2nd term in Eq. 1 is the new contribution from AGB stars. For this purpose we adopt the AGB yields of [16]. The 3rd term in Eq. 1 is the contribution from type Ia supernovae where m_{CO} is the mass the exploding carbon-oxygen white dwarf, and R_{Ia} is the SNIa supernova rate taken from [21]. Finally, the last two terms represent the trapping of elements in new stars, and the galactic infall rate (presumed to be of primordial material).

As noted by a number of authors, the yields of heavy magnesium isotopes in AGB stars is extremely temperature sensitive, and hence rather sensitive to detailed physics of the stellar models. Moreover, there are reasons to expect that the initial mass function at low metallicity could be biased toward intermediate-mass stars. One argument for this is simply that with fewer metals, the cooling is less efficient in the protostellar cloud, so that a more massive cloud is required to form a star. To account for these possibilities, we introduce a modest enhancement of the IMF for intermediate mass stars. Such an enhancement has often been proposed and is motivated by models for star formation at low metallicity. For example, it has been invoked [22] to account for a large population of white dwarfs as microlensing objects in the Galactic Halo. In fact, such a population of intermediate mass stars was recently proposed [23] to explain the dispersion of D/H observed in quasar absorption systems

[24]. Hence for the creation function, $B\Psi$, we write:

$$\begin{aligned} B(t)\Psi(m) &= B_1(t)\Psi_1(m) + B_2(t)\Psi_2(m) \\ &= B_1 m^{-2.35} + (B_2/m) \exp(-\log(m/5.0)^2/(2\sigma^2)). \end{aligned} \quad (2)$$

The IMF in Eq. (2) accounts for a standard Salpeter distribution of stellar masses, $\Psi_1(m)$, with the addition of a lognormal component of stars peaked at $5 M_\odot$, Ψ_2 . The dimensionless width, σ , is taken to be 0.07. This IMF is similar to a Gaussian with a width of $5\sigma = 0.35 M_\odot$. For the normal stellar component we take the time dependence as

$$B_1(t) = (1.0 - e^{-t/.5\text{Gyr}}) \sigma_{tot}(t) [\sigma_{gas}/\sigma_{tot}(t)]^2, \quad (3)$$

while for the intermediate mass component we take

$$B_2(t) = 5.5 e^{-t/.2\text{Gyr}} \sigma_{tot}(t) [\sigma_{gas}/\sigma_{tot}(t)]^2. \quad (4)$$

This is very similar to the model used in [23]. It contains an early burst of intermediate mass stars peaked at $5 M_\odot$ which is exponentially suppressed after 0.2 Gyr. The second component describes standard quiescent star formation with a smooth transition from the burst.

Figure 1 shows a comparison of our calculated magnesium isotope ratio vs iron abundance. The solid curve shows the result of the model described above including the AGB contribution. The QSO absorption-line systems in question have metallicities in the range from 0 to -2.5 with a typical iron abundance of $[\text{Fe}/\text{H}] \sim -1.5$. The mean isotopic ratio needed to account for the data of [1]-[4] is $^{25,26}\text{Mg}/^{24}\text{Mg} = 0.58$ (shown by the solid horizontal line) with a 1σ lower limit of 0.47 (dashed horizontal line). This figure clearly demonstrates that a plausible model is possible in which a sufficient abundance of heavy Mg isotopes can be produced to both explain the observed globular-cluster data and the apparent trends in the many-multiplet data or QSO absorption-line systems at high redshift.

The behavior in the evolution of the heavy isotopes can be explained as follows: Initially, the production of $^{25,26}\text{Mg}$ in the ejecta of intermediate mass stars is delayed by their relatively long lifetime (compared to very massive stars). Initial contributions to the chemical enrichment of the interstellar medium comes from the most massive and shortest lived stars. In this model, the burst of intermediate mass stars begins to produce $^{25,26}\text{Mg}$ at $[\text{Fe}/\text{H}] \gtrsim -2.5$. At this stage, during the intermediate mass burst, ^{25}Mg and ^{26}Mg are copiously produced relative to ^{24}Mg as per the yields of [16]. At higher metallicity, the ejecta from the standard population of (massive) stars which is poor in $^{25,26}\text{Mg}$ begins to dilute the ratio relative to ^{24}Mg , thereby producing the noticeable bump in $^{25,26}\text{Mg}/^{24}\text{Mg}$ around $[\text{Fe}/\text{H}] \sim -1.5$. At late times, the effect of the early generation of intermediate mass stars is largely washed away.

The dashed curve excludes the AGB yields and the intermediate mass component. It gives a result similar to that of [9]. We note that recently the AGB contribution was included in a chemical evolution model [25] for a

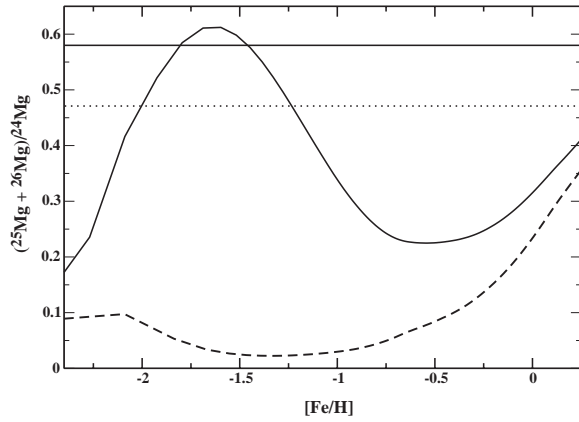


FIG. 1: The chemical evolution of the $^{25,26}\text{Mg}$ isotopes relative to ^{24}Mg . The solid curve is our result based on eq. (2) using the AGB Mg yields of [16]. The dashed curve is the result of turning off the AGB contribution and excluding the burst of intermediate mass stars. The horizontal lines indicate the ratio of $^{25,26}\text{Mg}/^{24}\text{Mg}$ necessary to explain the shifts seen in the many-multiplet analysis.

normal stellar distribution. While the results showed significantly higher abundances of $^{25,26}\text{Mg}$ relative to ^{24}Mg than that given by the dashed curve, they were not high enough to account for the claimed variability in α . An enhanced early population of intermediate mass stars is therefore necessary. Nevertheless, this seems more plau-

sible than a time varying fundamental constant.

We have argued that previous models for the apparent broadening of the Mg multiplet in QSO absorption-line systems may have left out the important possible contribution from the production of heavy magnesium isotopes during the AGB phase of low-metallicity intermediate-mass stars. We have shown that a simple, plausible galactic chemical evolution model can be constructed which explains both the large abundances of heavy Mg isotopes observed in globular clusters and the large abundance necessary to explain the many-multiplet data. We note that the hot bottom burning process in AGB stars is also likely to have altered the Si isotopes as well by proton captures on aluminum and silicon at the base of the convective envelope. Hence, it is possible that the supporting data from Si isotopes can be explained by this paradigm as well. Obviously more detailed work is warranted to clarify the ability of this mechanism to account for the data. Nevertheless, the model presented here is based upon plausible expectations of stellar and galactic evolution and should be taken seriously before demanding an alteration of any fundamental constant at high redshift.

We thank C. Cardall for helpful conversations. The work of K.A.O. was partially supported by DOE grant DE-FG02-94ER-40823. Work at the University of Notre Dame was supported by the U.S. Department of Energy under Nuclear Theory Grant DE-FG02-95-ER40934.

-
- [1] J. K. Webb, V. V. Flambaum, C. W. Churchill, M. J. Drinkwater and J. D. Barrow, *Phys. Rev. Lett.* **82** (1999) 884 [arXiv:astro-ph/9803165].
 - [2] M. T. Murphy *et al.*, *Mon. Not. Roy. Astron. Soc.* **327** (2001) 1208 [arXiv:astro-ph/0012419]. J. K. Webb *et al.*, *Phys. Rev. Lett.* **87** (2001) 091301 [arXiv:astro-ph/0012539].
 - [3] M. T. Murphy, J. K. Webb, V. V. Flambaum, C. W. Churchill and J. X. Prochaska, *Mon. Not. Roy. Astron. Soc.* **327** (2001) 1223 [arXiv:astro-ph/0012420].
 - [4] M. T. Murphy, J. K. Webb and V. V. Flambaum, arXiv:astro-ph/0306483.
 - [5] J. P. Uzan, *Rev. Mod. Phys.* **75** (2003) 403 [arXiv:hep-ph/0205340].
 - [6] V. A. Dzuba, V. V. Flambaum, and J. K. Webb, *Phys. Rev. A* **59** (1999) 230; *Phys. Rev. Lett.* **82** (1999) 888.
 - [7] J. N. Bahcall, C. L. Steinhardt and D. Schlegel, arXiv:astro-ph/0301507.
 - [8] P. Gay and D. L. Lambert, *Astrophys. J.* **533** (2000) 260 [arXiv:astro-ph/9911217].
 - [9] F. X. Timmes, S. E. Woosley and T. A. Weaver, *Astrophys. J. Suppl.* **98** (1995) 617 [arXiv:astro-ph/9411003].
 - [10] D. Yong, F. Grundahl, D. L. Lambert, P. E. Nissen and M. Shetrone, *Astron. Astrophys.* **402** (2003) 985 [arXiv:astro-ph/0303057].
 - [11] K.J.R. Rosman and P.D.P. Taylor, *J. Phys. Chem. Ref. Data*, **27** (1998) 1275.
 - [12] M.D. Shetrone, *Astron. J.* **112** (1996) 2639.
 - [13] D. Yong, D.L. Lambert, and I.I. Ivans, astro-ph/0309/079.
 - [14] M. Forestini and C. Charbonnel, *Astron. Astrophys. Suppl.* **123** (1996) 241.
 - [15] L. Siess, M. Livio, and J. Lattanzio, *Astrophys. J.* **570** (2002) 329 [arXiv:astro-ph/0201284].
 - [16] A.I. Karakas and J.C. Lattanzio, *PASA* **20** (2003) 279 and arXiv:astro-ph/0305011.
 - [17] S. E. Woosley and T. A. Weaver, *Astrophys. J. Suppl.* **101** (1995) 181.
 - [18] F.-K. Thielemann, K. Nomoto, and K. Yokoi, *Astron. Astrophys.* **158** (1986) 17.
 - [19] N. Iwamoto *et al.*, *Nucl. Phys. A* **719** (2003) 571.
 - [20] P. Marigo, *Astron. Astrophys.* **370** (2001) 194; L. Portinari, C. Chiosi, and A. Bressan, *Astron. Astrophys.* **334** (1998) 505.
 - [21] C. Kobayashi, T. Tsujimoto, and K. Nomoto, *Ap. J.* **539** (2000) 26.
 - [22] D. S. Ryu, K. A. Olive and J. Silk, *Astrophys. J.* **353** (1990) 81; B. D. Fields, G. J. Mathews and D. N. Schramm, *Astrophys. J.* **483** (1997) 625 [arXiv:astro-ph/9604095].
 - [23] B. D. Fields, K. A. Olive, J. Silk, M. Casse and E. Vangioni-Flam, *Ap. J.* **563** (2001) 653, [arXiv:astro-ph/0107389].
 - [24] D. Kirkman, D. Tytler, N. Suzuki, J. M. O'Meara and D. Lubin, arXiv:astro-ph/0302006.
 - [25] Y. Fenner *et al.*, arXiv:astro-ph/0307445.